

Thermal Buckling of Laminated Plates using Modified Mantari Function

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ABSTRACT

Critical buckling temperature of laminated plate under thermal load varied linearly along the thickness, is developed using a higher-order shape function which depends on a parameter ‘m’, which is improved to obtain results for thin and thick plates. Laminated plates’ equations of motion are obtained using virtual work principle and solved for simply supported boundary conditions. Angle and cross laminates thermal buckled mode shapes with different E_1/E_2 proportion, number of plies, (α_2/α_1) proportion, aspect ratios, are investigated. It is observed that this shape function gives thermal buckling for thin and thick plates but with $m = 0.05$ that agree well with other theories and linear distribution of temperature gives a rise to critical temperature approach to 50% than those caused by uniform thermal distribution.

Keywords: Linear Thermal Load; Non-uniform Thermal Load; Critical Temperature; Angle-ply Plate; Cross-ply Plate; Shear Deformation Theory

Introduction

Thermal buckling is an important failure mode for composite plate and shell when worked in thermal environment such as rockets, marine and space vehicles, aerospace structure and launch vehicles may be exposed to thermal loading due to the aerodynamic heating, as well as in other fields of modern engineering structures. Many researchers studied thermal buckling using different theories, therefore investigation of plate thermal buckling is of special interest.

Critical temperature of plate made from composite laminates is obtained using numerical methods such as differential quadrature method (DQM) [1]-

[2] and finite element method [3]-[7], these plates may be under uniform and/or non-uniform, linear or nonlinear thermal load with various boundary conditions, others [8]-[10] used variation methods to investigate buckling of laminated plates in thermal environment especially with different boundary conditions. Fazzolari and Carrera [11] and Fazzolari [12] carried out thermal stability and free vibration analyses of functionally graded (FG) and sandwich plates using Ritz and advanced Ritz method, taking into consideration uniform, linear and non-linear temperature rises through the-thickness layer plate. Xing and Wang [13], investigated thermal buckling for rectangular plates made from functionally graded material, under different boundary conditions using separation-of-variable method. Karami et al. [14], modified 3D plate theory to disperse wave in functionally graded nanoplates in hygro thermal environment and rest on an elastic foundation, Nicholas et al. [15] used genetic algorithm to maximize the critical temperature of the laminated composite plate subjected to different nonuniform thermal loads, therefore ply angle and stacking sequence are optimized.

Many researchers developed and used different higher-order shear deformation theory (HSDT) to study bending, buckling and response of plate with simply supported boundary [16]-[17], on a viscoelastic foundation or resting on a Winkler-Pasternak elastic foundation, in thermal or in a hygro-thermal environment or under hygro-thermo-mechanical load [18]-[20] which were as composite, functionally graded (FG), sandwich plate or made from porous functionally graded material [21]-[22], while effect of these environment on vibration behavior for these plates were investigated by [23]-[24]. Thermal stability of porous functionally graded material beam (of power law) under different thermal load distribution and boundary conditions were analyzed by [25]. Effect of reinforcement material (such as carbon nanotubes (CNTs)) on vibration behavior of microplates and shell in thermal environment were discussed by [26]-[27], using developed shear deformation theory with five or more unknown.

Different simple refined higher order shear deformation theory were developed and used to analyze thermo-mechanical buckling [28]-[31], bending and dynamic analyses of composite plates as in [32]-[34].

In present work, high order shear deformation theory is used to obtain critical temperature of simply supported laminated composite plate under linearly distributed temperature along the thickness, using displacement field proposed by Mantari et al. [35] but with $m = 0.05$ as developed by [36], and used it by [37] to obtain thermal buckling of laminated plates under uniform thermal load. Effect of many design parameters for thin and thick plate, E_1/E_2 ratio, aspect ratio, α_2/α_1 ratio for symmetric and antisymmetric cross and angle ply are investigated.

Displacements and strain

In present work, developed Mantari’s displacement field with $m = 0.05$ value is used to obtain critical thermal temperature of simply supported laminated plate, this modified field of displacement is, [36]:

$$\bar{u}(x, y, z) = u_x(x, y) + z \left(\frac{m\pi}{h} \theta_x - \frac{\partial w}{\partial x} \right) + \left(e^{m \cos \left(\frac{\pi z}{h} \right)} * \sin \frac{\pi z}{h} \right) * \theta_x \quad (1)$$

$$\bar{v}(x, y, z) = v_y(x, y) + z \left(\frac{m\pi}{h} \theta_y - \frac{\partial w}{\partial y} \right) + \left(e^{m \cos \left(\frac{\pi z}{h} \right)} * \sin \frac{\pi z}{h} \right) * \theta_y \quad (2)$$

$$\bar{w}(x, y, z) = w(x, y) \quad (3)$$

with: $m = 0.05$.

Linear strain based on displacement field are:

$$\epsilon_{xx} = \epsilon^0_{xx} + z\epsilon^1_{xx} + \sin \frac{\pi z}{h} e^{m \cos \left(\frac{\pi z}{h} \right)} \epsilon^2_{xx} \quad (4)$$

$$\epsilon_{yy} = \epsilon^0_{yy} + z\epsilon^1_{yy} + \sin \frac{\pi z}{h} e^{m \cos \left(\frac{\pi z}{h} \right)} \epsilon^2_{yy} \quad (5)$$

$$\gamma_{xy} = \epsilon^0_{xy} + z\epsilon^1_{xy} + \sin \frac{\pi z}{h} e^{m \cos \left(\frac{\pi z}{h} \right)} \epsilon^2_{xy} \quad (6)$$

$$\gamma_{xz} = \epsilon^0_{xz} + \left(-m * \sin^2 \left(\frac{\pi z}{h} \right) + \cos \frac{\pi z}{h} \right) \frac{\pi}{h} e^{m \cos \left(\frac{\pi z}{h} \right)} \epsilon^3_{xz} \quad (7)$$

$$\gamma_{yz} = \epsilon^0_{yz} + \left(-m * \sin^2 \left(\frac{\pi z}{h} \right) + \cos \frac{\pi z}{h} \right) \frac{\pi}{h} e^{m \cos \left(\frac{\pi z}{h} \right)} \epsilon^3_{yz} \quad (8)$$

where:

$$\left\{ \begin{matrix} \epsilon^0_{xx} \\ \epsilon^0_{yy} \\ \gamma^0_{xy} \end{matrix} \right\} = \left\{ \begin{matrix} \frac{\partial u_x}{\partial x} \\ \frac{\partial v_y}{\partial x} \\ \frac{\partial u_x}{\partial y} + \frac{\partial v_y}{\partial x} \end{matrix} \right\}, \left\{ \begin{matrix} \epsilon^1_{xx} \\ \epsilon^1_{yy} \\ \gamma^1_{xy} \end{matrix} \right\} = \left\{ \begin{matrix} \frac{m\pi}{h} \frac{\partial \theta_x}{\partial x} - \frac{\partial^2 w}{\partial x^2} \\ \frac{m\pi}{h} \frac{\partial \theta_y}{\partial y} - \frac{\partial^2 w}{\partial y^2} \\ \frac{m\pi}{h} \frac{\partial \theta_y}{\partial x} + \frac{m\pi}{h} \frac{\partial \theta_x}{\partial y} \\ -2 \frac{\partial^2 w}{\partial x \partial y} \end{matrix} \right\} \quad (9, 10)$$

$$\left\{ \begin{matrix} \epsilon^2_{xx} \\ \epsilon^2_{yy} \\ \gamma^2_{xy} \end{matrix} \right\} = \left\{ \begin{matrix} \frac{\partial \theta_x}{\partial x} \\ \frac{\partial \theta_y}{\partial y} \\ \frac{\partial \theta_y}{\partial x} + \frac{\partial \theta_x}{\partial y} \end{matrix} \right\} \quad (11)$$

$$\begin{Bmatrix} \Gamma_{xz}^0 \\ \gamma_{yz}^0 \end{Bmatrix} = \begin{Bmatrix} m \frac{\pi}{h} \Theta_x \\ m \frac{\pi}{h} \Theta_y \end{Bmatrix}, \quad \begin{Bmatrix} \gamma_{xz}^3 \\ \gamma_{yz}^3 \end{Bmatrix} = \begin{Bmatrix} \Theta_x \\ \Theta_y \end{Bmatrix} \quad (12 \text{ and } 13)$$

Equations of motion

To derive equations of motion, principle of virtual displacements is used with new shape function Reddy [38].

$$0 = \int_0^t \delta U + \delta V \quad (14)$$

For which:

$$\delta U = \int_{-\frac{h}{2}}^{\frac{h}{2}} \left\{ \int_0^k [\sigma_{xy} \delta \epsilon_{xy}^k + \sigma_{yz} \delta \epsilon_{yz}^k + \sigma_{xz} \delta \epsilon_{xz}^k + \sigma_{xx} \delta \epsilon_{xx}^k + \sigma_{yy} \delta \epsilon_{yy}^k] \partial x \partial y \right\} \partial z = 0 \quad (15)$$

$$\delta U = \int (N_1 \delta \epsilon_{xx}^0 + N_2 \delta \epsilon_{yy}^0 + N_6 \delta \epsilon_{xy}^0 + M_1 \delta \epsilon_{xx}^1 + M_2 \delta \epsilon_{yy}^1 + M_6 \delta \epsilon_{xy}^1 + P_1 \delta \epsilon_{xx}^2 + P_2 \delta \epsilon_{yy}^2 + P_6 \delta \epsilon_{xy}^2 + Q_2 \delta \epsilon_{yz}^0 + Q_1 \delta \epsilon_{xz}^0 + K_1 \delta \epsilon_{xz}^3 + K_2 \delta \epsilon_{yz}^3) \partial x \partial y = 0 \quad (16)$$

where, (N_i, M_i, P_i, Q_i and K_i) are the load results from the following integration:

$$\begin{aligned} (N_i, M_i, P_i) &= \sum_{k=1}^N \int_{z^{k-1}}^{z^k} \left(1, z, \sin\left(\frac{\pi z}{h}\right) * e^{m \cos\left(\frac{\pi z}{h}\right)} \right) * \sigma_i^k * dz \quad (i = 1, 2, 6) \\ (Q_1, K_1) &= \sum_{k=1}^N \int_{z^{k-1}}^{z^k} \left(1, \frac{\pi}{h} \left(-m * \sin^2\left(\frac{\pi z}{h}\right) + \cos\left(\frac{\pi z}{h}\right) \right) * e^{m \cos\left(\frac{\pi z}{h}\right)} \right) * \sigma_5^k * dz \\ (Q_2, K_2) &= \sum_{k=1}^N \int_{z^{k-1}}^{z^k} \left(1, \frac{\pi}{h} \left(-m * \sin^2\left(\frac{\pi z}{h}\right) + \cos\left(\frac{\pi z}{h}\right) \right) * e^{m \cos\left(\frac{\pi z}{h}\right)} \right) * \sigma_4^k * dz \end{aligned}$$

Substituting the virtual strain from Equations (4-13) into Equation (16) to get:

$$\begin{aligned} 0 = & - \int \left[\frac{\partial N_1}{\partial x} \delta u_x + \frac{\partial N_6}{\partial y} \delta u_x - \frac{\partial^2 M_1}{\partial x^2} \delta w + 2 \frac{\partial^2 M_6}{\partial x \partial y} \delta w + \frac{m\pi}{h} \frac{\partial M_2}{\partial y} \delta \theta_y - \right. \\ & \frac{\partial^2 M_2}{\partial y^2} \delta w + \frac{\partial N_2}{\partial y} \delta v_y + \frac{\partial N_6}{\partial x} \delta v_y + \frac{m\pi}{h} \frac{\partial M_6}{\partial y} \delta \theta_x + \frac{m\pi}{h} \frac{\partial M_1}{\partial x} \delta \theta_x + \frac{\partial P_6}{\partial y} \delta \theta_x + \\ & \frac{\partial P_1}{\partial x} \delta \theta_x - \frac{m\pi}{h} Q_1 \delta \theta_x - K_1 \delta \theta_x + \left(\frac{m\pi}{h} \frac{\partial M_6}{\partial x} \right) \delta \theta_y + \frac{\partial P_6}{\partial x} \delta \theta_y + \frac{\partial P_2}{\partial y} \delta \theta_y - \\ & \left. \left(\frac{m\pi}{h} Q_2 \right) \delta \theta_y - K_2 \delta \theta_y \right] \partial x \partial y \quad (17) \end{aligned}$$

Virtual work of thermal applied load δV is:

$$\delta V = \int_{\Omega} \left\{ N_x^T \delta \left(\frac{\partial w}{\partial x} \right)^2 + N_y^T \delta \left(\frac{\partial w}{\partial y} \right)^2 \right\} dx dy \quad (18)$$

Five equations of motion are derived by substituting Equations (17) and (18) into Equation (16):

$$\delta u: \frac{\partial N_{12}}{\partial y} + \frac{\partial N_1}{\partial x} = 0 \quad (19)$$

$$\delta v: \frac{\partial N_{12}}{\partial x} + \frac{\partial N_2}{\partial y} = 0 \quad (20)$$

$$\delta w: 2 \frac{\partial^2 M_{12}}{\partial x \partial y} + \frac{\partial^2 M_1}{\partial x^2} + \frac{\partial^2 M_2}{\partial y^2} + N_x^T \left(\frac{\partial^2 w}{\partial x^2} \right) + N_y^T \left(\frac{\partial^2 w}{\partial y^2} \right) = 0 \quad (21)$$

$$\delta \Theta_x: \frac{m\pi}{h} \frac{\partial M_1}{\partial x} + \frac{m\pi}{h} \frac{\partial M_{12}}{\partial y} + \frac{\partial P_1}{\partial x} + \frac{\partial P_{12}}{\partial y} - \frac{m\pi}{h} Q_1 - K_1 = 0 \quad (22)$$

$$\delta \Theta_y: \frac{m\pi}{h} \frac{\partial M_2}{\partial y} + \frac{m\pi}{h} \frac{\partial M_{12}}{\partial x} + \frac{\partial P_2}{\partial y} + \frac{\partial P_{12}}{\partial x} - \frac{m\pi}{h} Q_2 - K_2 = 0 \quad (23)$$

An orthotropic lamina stress-strain relation for a plane state of stress is [38]:

$$\begin{Bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{xy} \end{Bmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & Q_{16} \\ Q_{12} & Q_{22} & Q_{26} \\ Q_{16} & Q_{26} & Q_{66} \end{bmatrix} \begin{Bmatrix} \varepsilon_{xx} - \alpha_{xx} \Delta T \\ \varepsilon_{yy} - \alpha_{xx} \Delta T \\ \gamma_{xy} - 2\alpha_{xy} \Delta T \end{Bmatrix}, \begin{Bmatrix} \sigma_{yz} \\ \sigma_{xz} \end{Bmatrix} = \begin{bmatrix} Q_{44} & Q_{45} \\ Q_{45} & Q_{55} \end{bmatrix} \begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix} \quad (24)$$

The force results- strains related as:

$$\begin{Bmatrix} N_1 \\ N_2 \\ N_{12} \\ M_1 \\ M_2 \\ M_{12} \\ P_1 \\ P_2 \\ P_{12} \end{Bmatrix} = \begin{bmatrix} A_{11} & A_{12} & A_{16} & B_{11} & B_{12} & B_{16} & E_{11} & E_{12} & E_{16} \\ A_{12} & A_{22} & A_{26} & B_{12} & B_{22} & B_{26} & E_{12} & E_{22} & E_{26} \\ A_{16} & A_{26} & A_{66} & B_{16} & B_{26} & B_{66} & E_{16} & E_{26} & E_{66} \\ B_{11} & B_{12} & B_{16} & D_{11} & D_{12} & D_{16} & F_{11} & F_{12} & F_{16} \\ B_{12} & B_{22} & B_{26} & D_{12} & D_{22} & D_{26} & F_{12} & F_{22} & F_{26} \\ B_{16} & B_{26} & B_{66} & D_{16} & D_{26} & D_{66} & F_{16} & F_{26} & F_{66} \\ E_{11} & E_{12} & E_{16} & F_{11} & F_{12} & F_{16} & H_{11} & H_{12} & H_{16} \\ E_{12} & E_{22} & E_{26} & F_{12} & F_{22} & F_{26} & H_{12} & H_{22} & H_{26} \\ E_{16} & E_{26} & E_{66} & F_{16} & F_{26} & F_{66} & H_{16} & H_{26} & H_{66} \end{bmatrix} \begin{Bmatrix} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \varepsilon_{xy}^0 \\ \varepsilon_x^1 \\ \varepsilon_y^1 \\ \varepsilon_{xy}^1 \\ \varepsilon_x^2 \\ \varepsilon_y^2 \\ \varepsilon_{xy}^2 \end{Bmatrix} \quad (25)$$

$$\begin{pmatrix} Q_1 \\ Q_2 \\ K_1 \\ K_2 \end{pmatrix} = \begin{bmatrix} A_{44} & A_{45} J_{44} & J_{45} \\ A_{45} & A_{55} J_{45} & J_{55} \\ J_{44} & J_{45} L_{44} & L_{45} \\ J_{45} & J_{55} L_{45} & L_{55} \end{bmatrix} \begin{pmatrix} \gamma_{yz}^0 \\ \gamma_{xz}^0 \\ \gamma_{yz}^3 \\ \gamma_{xz}^3 \end{pmatrix} \quad (26)$$

$$\begin{pmatrix} N_x^T \\ N_y^T \end{pmatrix} = \sum_{k=1}^N \int_{z^k}^{z^{k+1}} \begin{bmatrix} Q_{11} & Q_{12} & Q_{16} \\ Q_{12} & Q_{22} & Q_{26} \end{bmatrix} \begin{pmatrix} \alpha_{xx} \\ \alpha_{yy} \\ 2\alpha_{xy} \end{pmatrix} \Delta T dz \quad (27)$$

$$\begin{pmatrix} M_x^T \\ M_y^T \end{pmatrix} = \sum_{k=1}^N \int_{z^k}^{z^{k+1}} \begin{bmatrix} Q_{11} & Q_{12} & Q_{16} \\ Q_{12} & Q_{22} & Q_{26} \end{bmatrix} \begin{pmatrix} \alpha_{xx} \\ \alpha_{yy} \\ 2\alpha_{xy} \end{pmatrix} \Delta T z dz \quad (28)$$

where:

$$(A_{ij}, B_{ij}, D_{ij}, E_{ij}) = \int_{-\frac{h}{2}}^{\frac{h}{2}} Q_{ij}(1, z, z^2, \sin(\frac{\pi z}{h})) e^{m \cos(\frac{\pi z}{h})} dz \quad (i, j=1, 2, 6) \quad (29)$$

$$(F_{ij}, H_{ij}) = \int_{-\frac{h}{2}}^{\frac{h}{2}} Q_{ij}(z * \sin(\frac{\pi z}{h})) e^{m \cos(\frac{\pi z}{h})}, e^{2m \cos(\frac{\pi z}{h})} * \sin^2(\frac{\pi z}{h}) dz \quad (i, j=1, 2, 6) \quad (30)$$

$$J_{ij} = \int_{-\frac{h}{2}}^{\frac{h}{2}} Q_{ij} \frac{\pi}{h} e^{m \cos(\frac{\pi z}{h})} (-m * \sin^2(\frac{\pi z}{h}) + \cos(\frac{\pi z}{h})) * dz \quad (31)$$

$$L_{ij} = \int_{-\frac{h}{2}}^{\frac{h}{2}} Q_{ij} (\frac{\pi}{h})^2 e^{2m \cos(\frac{\pi z}{h})} (-m * \sin^2(\frac{\pi z}{h}) + \cos(\frac{\pi z}{h}))^2 dz \quad (i, j = 4, 5) \quad (32)$$

while $\Delta T = \frac{\Delta T}{h} (z + \frac{h}{2})$ and Q_{ij} ($i, j=1, 2, 4, 5, 6$) = plane stress reduced stiffness.

Navier solution

Navier's series for simply supported cross laminates, are [38]:

$$u_x(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{U}_{mn}(t) * \cos(\frac{m\pi}{a} x) * \sin(\frac{n\pi}{b} y) \quad (33)$$

$$v_y(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{V}_{mn}(t) * \sin(\frac{m\pi}{a} x) * \cos(\frac{n\pi}{b} y) \quad (34)$$

$$w(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{W}_{mn}(t) * \sin(\frac{m\pi}{a} x) * \sin(\frac{n\pi}{b} y) \quad (35)$$

$$\theta_x(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{\theta}_{x_{mn}}(t) * \cos\left(\frac{m\pi}{a}x\right) * \sin\left(\frac{n\pi}{b}y\right) \quad (36)$$

$$\theta_y(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{\theta}_{y_{mn}}(t) * \sin\left(\frac{m\pi}{a}x\right) * \cos\left(\frac{n\pi}{b}y\right) \quad (37)$$

While for simply supported Angle-Ply laminates, the following series are used [37]:

$$u(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{U}_{mn}(t) * \sin\left(\frac{m\pi}{a}x\right) * \cos\left(\frac{n\pi}{b}y\right) \quad (38)$$

$$v(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{V}_{mn}(t) * \cos\left(\frac{m\pi}{a}x\right) * \sin\left(\frac{n\pi}{b}y\right) \quad (39)$$

$$w(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{W}_{mn}(t) * \sin\left(\frac{m\pi}{a}x\right) * \sin\left(\frac{n\pi}{b}y\right) \quad (40)$$

$$\theta_x(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{\theta}_{x_{mn}}(t) * \sin\left(\frac{m\pi}{a}x\right) * \cos\left(\frac{n\pi}{b}y\right) \quad (41)$$

$$\theta_y(x, y, t) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \bar{\theta}_{y_{mn}}(t) * \cos\left(\frac{m\pi}{a}x\right) * \sin\left(\frac{n\pi}{b}y\right) \quad (42)$$

where: a, b are dimensions of plate.

Eignvalue problem

Equations of motion as function of displacement componants are obtained by substituting force and moment resultants from Equations (25 - 28) in Equations (19-23) and then by substituting Equations (33-37) for cross ply and Equations (38-42) for angle ply, the following eignvalue problem is obtained:

$$\begin{pmatrix} k_{11} & k_{12} & & k_{13} & & k_{14} & k_{15} \\ & k_{22} & & k_{23} & & k_{24} & k_{25} \\ & & k_{33} - \alpha^2(N_x^T + M_x^T) - \beta^2(N_y^T + M_y^T) - \alpha\beta N_{xy}^T & & & k_{34} & k_{35} \\ & & & & & k_{44} & k_{45} \\ & & & & & & k_{55} \end{pmatrix} \{e\} = 0 \quad (43)$$

where: $\{e_{ij}\} = \{U_{mn}, V_{mn}, W_{mn}, \theta_{x_{mn}}, \theta_{y_{mn}}\}$, and k_{ij} is element of stiffness, from Equation (43) thermal buckling of the plate is calculated.

Results and Discussion

Critical temperature for angle and cross laminated plates under non-uniform temperature distribution is obtained by programming present solution using MATLAB15. To verify the suggested above solution, the obtained results compared with other works as shown in Table 1 and Table 2, which agree well. The effect of design parameters on critical temperature of laminated plates are analyzed such as, symmetric and antisymmetric ply, number of layers, orthotropy ratio E_1/E_2 , thermal expansion coefficient ratio (α_2/α_1) under uniform and linear temperature distribution.

Table 1: Normalized critical temperature for isotropic square plate, $E_1 = E_2 = 380 \text{ Gpa}$, $\nu = 0.3$, $\alpha = 7.4e-6$

a/H	100	80	60	40	20	10
Present	31.15	48.67	86.55	194.479	771.269	2961.0484
Cetkovic [4]	34.182	53.395	94.871	213.113	845.027	3266.311
CLPT	24.198	43.434	84.995	203.738	844.955	3409.821
HODT	24.177	43.387	84.484	202.984	833.032	3224.968

Table 2: Critical temperature ($T_{cr} = T^*1e-6$) for isotropic square plate ($a/H = 100$), $E = 1$, $\nu = 0.3$, $\alpha = 2e-6$

Tcr			
Present work	Cetkovic [4]	Kari et al. [39]	Discrepancy%
115.256	126.5	126	9.5

Critical buckling temperature of thick plate is larger than those for thin plate while uniform temperature distribution cause plate buckle under smaller temperature than those under linear distribution as shown in Table 3. Symmetric and antisymmetric cross ply $[0/90]_2$ and $[0/90]_4$ for square plate with different thickness ratio, are obtained in Table 4, as expected that critical temperature increases as number of layer increases, but it decreases with increasing thicknes ratio and it is larger for antisymmetric, for following tables and figures the layers properties used are: $E_1/E_2 = 25$, $G_{12} = G_{13} = 0.5 E_2$, $G_{23} = 0.2 E_2$, $\nu_{12} = \nu_{13} = \nu_{23} = 0.25$, $\alpha_2/\alpha_1 = 3$.

Normalized critical temperature for atisymmetric angle ply thick and thin square plate are listed in Table 5, which shows that increasing (a/H) ratio decreases critical temperature. Different critical thermal buckling modes for plates for $[30/-30]_4$ angle-ply square plate are shown in Figures 1 and 2. For these four figures the dimensionless critical temperature is ($T^*_{cr} = T^*a^2h/\pi^2D_{22}$).

Table 3: Critical temperature T_{cr} for cross-ply square plate

a/H	Tcr		Discrepancy %
	Uniform	Linear	
4	0.0109699	0.022893	52
10	0.002874	0.005846	50.8
20	0.0007931	0.00159957	50.4

Table 4: Effect of lamine type on critical temperature T_{cr} for cross-ply square plate

a/H	Symmetric cross-ply		Anti-Symmetric cross-ply	
	(0/90) ₂	(0/90) ₄	(0/90) ₂	(0/90) ₄
5	0.0197	0.02431	0.02327	0.02501
10	0.009699	0.01120	0.01018	0.01124
50	0.0006238	0.0006303	0.00053968	0.0006088

Table 5 Normalized critical temperature [$T_{cr} = T^* \alpha_1 * 10^*(b/H)^2$] for angle-ply square plate

a/H	Tcr		
	[30/-30] ₄	[45/-45] ₄	[60/-60] ₄
5	7.3041	7.5790	7.3041
10	15.4472	16.5984	15.4472
50	24.5405	27.5317	24.5405

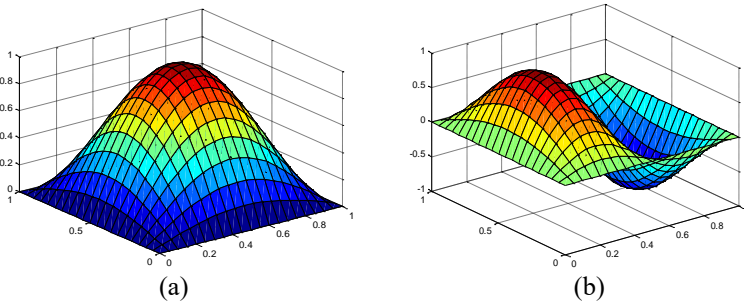


Figure 1: Normalized thermal buckling mode of angle-ply $[\pm 30]_4$ square plate, a) mode shape ($m = 1, n = 1$), b) mode shape ($m = 2, n = 1$)

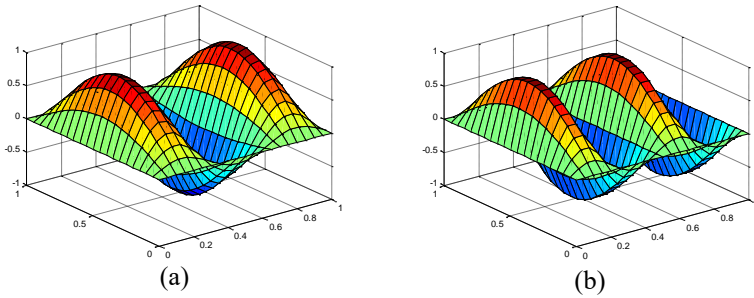


Figure 2: Normalized thermal buckling mode of angle-ply $[\pm 30]_4$ square plate, a) mode shape ($m = 3, n = 1$), b) mode shape ($m = 4, n = 1$)

Effect of changing (E_1/E_2) on critical temperature for eight layers anti symmetric angle ply plates for different angle ply are shown in Figure 3, as expected when increasing orthotropy ratio will increases buckling temprature (it seems reversly because T_{cr} is divided by D_{22}), also effect of thermal expansion coefficient ratio (α_2/α_1) on critical buckling temperature of eight layer laminated, are obtained in Figure 4 as expected critical temperature decreases when (α_2/α_1) increase since it decreases stiffness of plate, and angle plies plate has thermal buckling behavior better than cross plate, thickness and aspect ratio are ($a/H = 10, a/b = 1$).

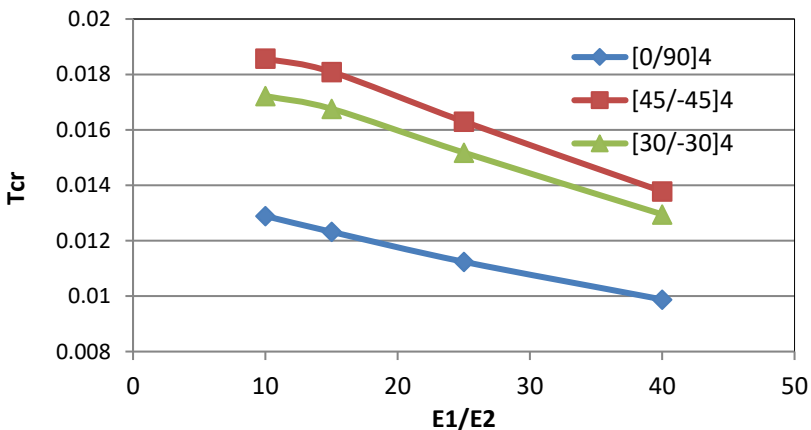


Figure 3: Normalized thermal buckling mode for antisymmetric plate, mode shape ($m = 1, n = 1$)

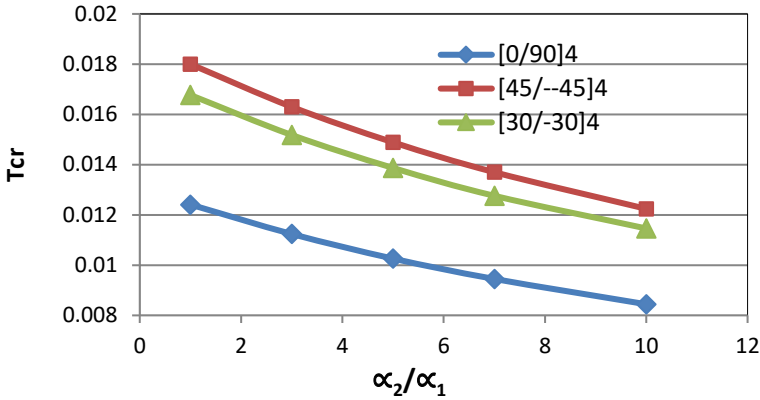


Figure 4: Normalized thermal buckling mode for antisymmetric plate, mode shape ($m = 1, n = 1$)

Maximum critical buckling temperature is for plates with (45°) angle as listed in Table 6, and it is greater for plate under linear than uniform temperature distribution.

Table 6: Normalized critical temperature [$T_{cr} = T^* \alpha_1 \cdot 10^*(b/H)^2$] for angle-square plate

Angle	No. of layers	Tcr		Discrepancy %
		Linear	Uniform [40]	
30	2	9.6671	4.8461	49.8
	4	19.3646	9.5515	50.6
	8	21.7228	10.6972	50.7
45	2	10.0923	5.0503	48
	4	21.3294	10.5094	50.7
	8	24.0398	11.8277	50.7
60	2	9.6671	4.8461	49.8
	4	19.3646	9.5515	50.6
	8	21.7228	10.6972	50.7

Conclusions

Mantari displacement field but with ‘ $m = 0.05$ ’ developed by Majeed and Sadiq [36], is used to analyze critical buckling temperature of cross and angle laminated composite thick and thin plates, under linear distribution of temperature and gives results agree well with those obtained by other theories. As expected critical temperature decreases as thickness ratio and thermal

expansion coefficient ratio (α_2/α_1) increase while it increases with the increase of aspect ratio and orthotropy ratio, and angle plies plate has thermal buckling behavior better than cross plate, while critical temperature of laminated plates under linear distribution of temperature along thickness is higher than that under uniform thermal distribution approach (50%), also the thermal buckling mode of these plates are not change.

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